VERY LOW MASS STELLAR AND SUBSTELLAR COMPANIONS TO SOLAR-LIKE STARS FROM MARVELS.
IV. A CANDIDATE BROWN DWARF OR LOW-MASS STELLAR COMPANION TO HIP 67526

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ABSTRACT

We report the discovery of a candidate brown dwarf (BD) or a very low mass stellar companion (MARVELS-5b) to the star HIP 67526 from the Multi-object Apache point observatory Radial Velocity Exoplanet Large-area Survey (MARVELS). The radial velocity curve for this object contains 31 epochs spread over 2.5 yr. Our Keplerian fit, using a Markov Chain Monte Carlo approach, reveals that the companion has an orbital period of 90.2695±0.0188 days, an eccentricity of 0.4375 ± 0.0040, and a semi-amplitude of 2948.14±16.65 m s⁻¹. Using additional high-resolution spectroscopy, we find the host star has an effective temperature $T_{\text{eff}} = 6004 \pm 34$ K, a surface gravity log $g$ (cgs) = 4.55 ± 0.17, and a metallicity [Fe/H] = +0.04 ± 0.06. The stellar mass and radius determined through the empirical relationship of Torres et al. yields $1.10 \pm 0.09 M_\odot$ and $0.92 \pm 0.19 R_\odot$. The minimum mass of MARVELS-5b is $65.0 \pm 2.9 M_{\text{Jup}}$, indicating that it is likely to be either a BD or a very low mass star, thus occupying a relatively sparsely populated region of the mass function of companions to solar-type stars. The distance to this system is 101 ± 10 pc from the astrometric measurements of Hipparcos. No stellar tertiary is detected in the high-contrast images taken by either FastCam lucky imaging or Keck adaptive optics imaging, ruling out any star with mass greater than 0.2 $M_\odot$ at a separation larger than 40 AU.

Key words: binaries: spectroscopic – brown dwarfs – stars: individual (HIP 67526) – stars: low-mass – techniques: radial velocities

Online-only material: color figures

1. INTRODUCTION

Brown dwarfs (BDs; Basri 2000) are star-like objects that are not massive enough to sustain stable hydrogen burning but are sufficiently massive to fuse deuterium (Chabrier et al. 2000; Spiegel et al. 2011). As a result, their luminosity and temperature drop throughout their lifetimes (e.g., Burrows et al. 1997; Baraffe et al. 2003). To date, over 800 BDs have been directly and indirectly discovered through a variety of methods (e.g., Rebolão et al. 1995; Oppenheimer et al. 1995; Ruiz et al. 1997; 24 LAMOST Fellow.)
In this paper, we report a candidate BD or a low-mass stellar companion (MARVELS-5b) to HIP 67526 with a period of ~90 days from MARVELS. In Section 2.1, we describe the RV measurements and solve for the spectroscopic orbital elements using Markov Chain Monte Carlo (MCMC) analysis. We analyze the photometric data from SuperWASP and the astrometric data from Hipparcos in Sections 2.2 and 2.3, respectively. In Section 3.1, we determine precise stellar parameters for the primary star. Using the stellar mass derived in Section 3.2, we then estimate the mass of the companion in Section 3.3. The evolutionary state of the host star is studied in Section 3.4. The high-contrast imaging is presented in Section 3.5. Finally, we provide a discussion and a summary in Section 4.

2. OBSERVATIONS AND RESULTS FOR THE LOW-MASS COMPANION

2.1. Differential Radial Velocities

2.1.1. MARVELS and TNG/SARG Measurements

HIP 67526 was selected as an RV survey target according to the MARVELS preselection criterion (Lee et al. 2011). It has been monitored at 21 epochs using the MARVELS

http://www.sdss3.org/surveys/marvels.php
HIP 67526 was identified as a star bearing an unseen companion by performing Lomb–Scargle (L-S) periodogram analysis (e.g., Lomb 1976; Scargle 1982; Cumming 2004; Balvev 2008) on the 21 MARVELS RV points. There are two significant peaks on the L-S periodogram with periods at \(~88\) days and \(~46\) days (Figure 1). The false alarm probability (hereafter FAP) of the 88 day peak is 0.00367%, and the FAP of the 46 day peak is 0.0201%. We fit a Keplerian orbit to the observed RV curve, forcing the period to be close to \(~88\) days and \(~46\) days. The preliminary fitting results are illustrated in Figure 2. The solution at an orbital period of 90.2 days provides a better fit to the MARVELS RV curve than the solution at an orbital period of 45.6 days. The shorter orbital period peak in the periodogram is probably an alias. The minimum mass (if \(\sin i = 1\)) of the unseen companion from the longer period solution is \(\sim 65 M_{\text{Jup}}\) (see Section 3.3 for details). The estimated minimum mass is below the hydrogen burning limit and places MARVELS-5b within the sparsely populated region of the mass function of companions to solar-like stars.

We collected 10 additional RV measurements with the SARG spectrograph (Gratton et al. 2001) at the 3.58 m Telescopio Nazionale Galileo (TNG) Telescope in late 2010 and 2011. The spectrograph covers a wavelength range of 4620–7920 Å with \(R \sim 57,000\). The simultaneous iodine cell technique (Butler et al. 1996) was employed to calibrate the RV measurements. The raw spectra were reduced by using the standard IRAF\(^{26}\) Echelle reduction packages. The final extracted differential RVs from MARVELS and TNG/SARG are presented in Table 1. The RV curve was sampled in total at 31 epochs using these two instruments over 2.5 yr.

### 2.1.2. Spectroscopic Orbital Elements

We have performed a Bayesian analysis of the observed RVs using a model consisting of the primary star and one low-mass companion on an eccentric Keplerian orbit based on

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\(^{26}\) http://iraf.noao.edu/
the combined differential RV observations of MARVELS and TNG/SARG.

We calculated a posterior sample using the MCMC technique as described in Ford (2006). Each state in the Markov Chain is described by the parameter set \( \theta = \{ P, K, e, \omega, M, \gamma_M, \gamma_T, \sigma_j \} \), where \( P \) is the orbital period, \( K \) is the velocity semi-amplitude, \( e \) is the orbital eccentricity, \( \omega \) is the argument of periastron, and \( M \) is the mean anomaly at the chosen epoch \( t_r \). The parameters \( \gamma_M \) and \( \gamma_T \) are constant systemic velocity terms for the MARVELS and TNG/SARG instruments, respectively, used to account for the offsets between the observed differential RV data and the zero point of the Keplerian RV model. The “jitter” parameter, \( \sigma_j \), describes any excess noise (Wright 2005), including both astrophysical sources of noise (e.g., stellar oscillation, stellar spots) and any instrumental noise not accounted for in the quoted measurement uncertainties. We use standard priors for each parameter (see Ford & Gregory 2007). The prior is uniform in the log of the orbital period \( P \), while for \( K \) and \( \sigma_j \) we used a modified Jeffreys’s prior (Gregory 2005). Priors for the remaining parameters are uniform: \( e \) (between zero and unity), \( \omega \) and \( M \) (between zero and \( 2\pi \)), \( \gamma_M \) and \( \gamma_T \). Following Ford (2006), we adopt a likelihood (i.e., conditional probability of making the specified measurements given a particular set of model parameters) of

\[
p(v | \theta, M) \propto \prod_k \exp \left\{ -\frac{(v_k,0 - v_k,\theta)^2}{2\sigma_k^2} \right\},
\]

where \( v_k \) is observed velocity at time \( t_k \), \( v_k,\theta \) is the model velocity at time \( t_k \) given the model parameters \( \theta \), and \( \sigma_k \) is the measurement uncertainty for the observation at time \( t_k \).

To test the robustness of the MCMC analysis, we calculate five Markov Chains starting from different initial states, each for \( 5 \times 10^7 \) steps. To prevent the choice of initial states from influencing our results, we consider only the second half of each chain. We calculate the Gelman–Rubin test statistic (which compares the variance of a parameter within each chain to the variance between chains; Gelman & Rubin 1992) for each model parameter. We find no indications that the Markov Chains have yet to converge and conclude that the Markov Chains provide an adequate posterior sample for inferring the orbital parameters and uncertainties.

We combine the Markov Chains described above to estimate the joint posterior probability distribution for the orbital model of HIP 67526. For orbital eccentricity, we also used the \( \Gamma \) method described in Wang & Ford (2011), which leads to a result similar to that from the MCMC analysis. The median values are taken for each model parameter based on the marginal posterior probability distributions. The uncertainties are calculated as the standard deviation about the mean value from the combined posterior sample. Since the shape of the marginal posterior distribution is roughly similar to a multivariate normal distribution, the median value plus or minus the reported uncertainty roughly corresponds to a 68.3% confidence interval. Finally, we convert the model parameters to traditional standard parameters of a spectroscopic orbit and report the results in Table 2. The phase-folded RV curve is presented in Figure 3.

### 2.2. SuperWASP Photometry

We searched the SuperWASP public archived database (Butters et al. 2010) and found 1378 photometric data measurements of HIP 67526 observed in 2004 and 5680 data points in 2007. The mean absolute deviation of the light curve is 9.7 mmag. We first searched for a transit-like dip in brightness at short periods between 0.2 and 10 days. We find no significant detection of a transit event. Next, we searched for transits specifically in the range of 85–95 days, which includes the best-fit period from the spectroscopic RV curve. The phase-folded data are sparsely covered at these long periods, and we find no significant transit signal. In summary, we do not find a transit in SuperWASP photometric data with a long or a short period. We also attempted to search for a sinusoidal signal in the light curve but found no significant signal.

### 2.3. Hipparcos Astrometry

HIP 67526 exists in the Hipparcos catalog with a parallax distance of 100 ± 10 pc from the Sun. It is possible that the orbital motion of the star due to the gravitational influence of

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<table>
<thead>
<tr>
<th>Table 1</th>
<th>Differential Radial Velocity Measurements</th>
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<td>Instrument</td>
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Note. \(^{a}\) Heliocentric Julian Day.

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<th>Table 2</th>
<th>Orbital Elements of MARVELS-5b</th>
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<td>Parameter</td>
<td>Units</td>
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<tr>
<td>( P )</td>
<td>Period (days)</td>
</tr>
<tr>
<td>( K )</td>
<td>RV semi-amplitude (m s(^{-1}))</td>
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<tr>
<td>( e )</td>
<td>Eccentricity</td>
</tr>
<tr>
<td>( \omega )</td>
<td>Argument of periastron (deg)</td>
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<tr>
<td>( T_0 )</td>
<td>Epoch of periastron (HJD)</td>
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<tr>
<td>( \gamma_M )</td>
<td>MARVELS systemic velocity (m s(^{-1}))</td>
</tr>
<tr>
<td>( \gamma_T )</td>
<td>TNG/SARG systemic velocity (m s(^{-1}))</td>
</tr>
<tr>
<td>( \sigma_j )</td>
<td>Jitter (m s(^{-1}))</td>
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where \( v_k \) is observed velocity at time \( t_k \), \( v_k,\theta \) is the model velocity at time \( t_k \) given the model parameters \( \theta \), and \( \sigma_k \) is the measurement uncertainty for the observation at time \( t_k \).
its companion can be resolved by Hipparcos astrometry. This would allow the inclination \( i \) and the ascending node \( \Omega \) of the Keplerian orbit, and thus the true mass of MARVELS-5b, to be well constrained (Sahlmann et al. 2011). We retrieved the dataset of HIP 67526 from the Intermediate Astrometric Data (IAD) of the new Hipparcos reduction (van Leeuwen 2007), including the satellite orbit number, the epoch \( t \), the parallax factor \( \Pi \), the scan angle orientation \( \psi \), the absissa residual \( \delta A \), and the absissa error \( \sigma_\delta \) for every satellite scan. There are 123 available Hipparcos scans on HIP 67526 in the IAD and the average absissa error is \( \bar{\sigma}_\delta \sim 10 \) mas. Thus, the dataset allows a 1\( \sigma \) detection of an orbit with an angular size of \( \sigma_\delta/\sqrt{N} = 10/\sqrt{123} \sim 1 \) mas.

We then estimate the minimum angular semimajor axis (in mas) of the primary’s orbit, which can be written as

\[
\alpha_s \sin i = 3.35729138 \times 10^{-5} K P \sqrt{1 - e^2} \sigma, \tag{2}
\]

where \( K \) (in m s\(^{-1}\)), \( P \) (in yr), and \( e \) are the spectroscopic orbital elements, \( \sigma \) (in mas) is the parallax, and \( i \) is the unknown inclination (Pourbaix 2001). This equation yields a minimum angular semimajor axis \( \sim 0.2 \) mas for HIP 67526. Therefore, for nearly edge-on orbits, the angular size of the primary’s orbit is well below the 1\( \sigma \) detection threshold, and thus the motion of HIP 67526 about the system’s center-of-mass cannot be detected for such geometries. Assuming that the Hipparcos data of HIP 67526 are consistent with no astrometric signal from the orbit around the center-of-mass of the system, and that orbits of \( \sim 1 \), \( \sim 2 \), and \( \sim 3 \) mas would have been detected at 1\( \sigma \), 2\( \sigma \), and 3\( \sigma \), we can place upper limits on the companion mass of \( \sim 0.33 \, M_\odot \) (1\( \sigma \)), and \( \sim 0.80 \, M_\odot \) (2\( \sigma \)), and \( 1.49 \, M_\odot \) (3\( \sigma \)). As argued in Section 3.3, such massive companions are a posteriori unlikely to be flat or falling priors on the companion mass distribution. For priors that increase with increasing mass, companions of mass \( \geq 0.5 \, M_\odot \) are not a posteriori implausible, but would be ruled out based on the lack of evidence of a second set of spectral lines in the high-resolution spectra, if the companion was luminous (i.e., not a remnant).

3. OBSERVATIONS AND RESULTS FOR THE HOST STAR

3.1. Spectroscopic Parameters and Spectral Energy Distribution Analysis

In order to characterize the host star HIP 67526, two moderate-resolution spectra \((R \sim 31,500)\) were taken with the ARC Echelle Spectrograph (ARCES; Wang et al. 2003) mounted on the Apache Point Observatory 3.5 m telescope on UT 2010 June 10. The spectra cover the full optical range from 3600 Å to 1.0 \( \mu \)m. The spectra were obtained using the default 1.6 × 3.2 slit and an exposure time of 1200 s. The raw data were processed using standard IRAF techniques. The extracted one-dimensional spectra were converted to vacuum wavelengths and to the heliocentric frame. The data were normalized by fitting a series of polynomials to the continuum. We utilized two individual pipelines to derive basic stellar parameters such as \( T_{\text{eff}} \), \( \log g \), and \([\text{Fe/H}]\) for the host star. Both pipelines are based on the requirements of excitation and ionization equilibria of Fe\( i \) and Fe\( ii \). However, different versions of ATLAS9 plane-parallel model atmospheres (Kurucz 1993 and Castelli & Kurucz 2004) and different iteration algorithms are implemented. We refer the readers to Wisniewski et al. (2012) for more details on the pipelines. The derived stellar parameters from these two pipelines are usually consistent to within 1\( \sigma \) of the associated errors. Thus, we simply adopted the weighted average values as the final determined stellar parameters. We combined the internal errors from the two pipelines as \( 1/\sigma^2 = 1/\sigma_1^2 + 1/\sigma_2^2 \) for each parameter, and
added in quadrature a systematic error of 18 K, 0.08, 0.03, and 0.02 km s\(^{-1}\) for \(T_{\text{eff}}, \log g, [\text{Fe}/H], \) and \(V_{\text{mic}}\), respectively (Wisniewski et al. 2012). The final results are summarized in Table 3.

We collected the optical and near-infrared absolute photometry of HIP 67526 from the \textit{Hipparcos}, Two Micron All Sky Survey (2MASS), and \textit{Wide-field Infrared Survey Explorer} (WISE) catalogs (Table 3) to construct a spectral energy distribution (SED; see Figure 4) and fit with a NextGen model atmosphere (Hauschildt et al. 1999). The resultant stellar parameters, \(T_{\text{eff}} = 5800 \pm 200\) K, \(\log g\) (cgs) = 4.0 \(\pm\) 1.0 and \([\text{Fe}/H] = 0.0 \pm 0.5\), are in good agreement with the parameters derived from the SED within the errorbars. In addition, the SED fitting indicates that HIP 67526 suffers only slight extinction \(A_V = 0.035 \pm 0.035\).

### 3.2. Stellar Mass and Radius

We determine the stellar mass and radius using two methods. First, we use the empirical relationship of Torres et al. (2010) with our values for \(T_{\text{eff}}, \log g,\) and \([\text{Fe}/H]\). Uncertainties in the mass and radius are derived by adding in quadrature the correlations of the best-fit coefficients from Torres et al. (2010) and the scatter in the relation as reported in their study. The correlations between the stellar parameters \(T_{\text{eff}}, \log g,\) and \([\text{Fe}/H]\) are not measured and are therefore not considered. We find a mass \(M_* = 1.10 \pm 0.09 M_\odot\) and a radius \(R_* = 0.92 \pm 0.19 R_\odot\).

The existence of a trigonometric parallax provides additional information to constrain the mass and radius of the primary star. We incorporate this data by running an MCMC analysis that fully explores parameter space. One million iterations in the MCMC were run, stepping through \(T_{\text{eff}}, \log g, [\text{Fe}/H],\) parallax (\(\sigma\)), and \(A_V\). We use random starting values to initiate the chain. For each iteration, we calculate a mass and radius following Torres et al. (2010) and the iteration’s values of \(T_{\text{eff}}, \log g,\) and \([\text{Fe}/H]\). A stellar luminosity is calculated via the Stefan–Boltzmann law, then a bolometric correction to the 2MASS \(K\) band is applied by interpolating the table of corrections as a function of \(T_{\text{eff}}\) for \([M/H] = 0.0\) and \(g = 4.5\) from Masana et al. (2006). The absolute \(K\) magnitude is calculated from the luminosity and bolometric correction, after which the apparent magnitude is calculated from the absolute magnitude and the iteration’s values of \(\sigma\) and \(A_V\).

After each iteration, a \(\chi^2\) statistic is calculated as the sum of the individual \(\chi^2\) for \(T_{\text{eff}}, \log g, [\text{Fe}/H], \sigma,\) and \(A_V\), where the expected values for \(T_{\text{eff}}, \log g,\) and \([\text{Fe}/H]\) are the values determined spectroscopically, the expected value for \(\sigma\) comes from the \textit{Hipparcos} catalog, and the expected value for \(A_V\) comes from the SED analysis. The next iteration’s trial parameters are selected using Gaussians centered on the current iteration’s values with widths equal to the 1σ parameter uncertainties for \(T_{\text{eff}}, [\text{Fe}/H],\) and \(A_V\), and 0.1σ for \(\log g\) and \(\sigma\). These widths were empirically determined such that the overall trial acceptance rate was \(\sim 24\%\), close to the optimal value for multi-dimensional chains (Gelman et al. 2003).

The first 1% of iterations are rejected as a burn-in period, while the remaining iterations are used to determine the best-fit final parameters \((M_*, R_*, T_{\text{eff}}, \log g, [\text{Fe}/H], \sigma, A_V)\). The 1σ uncertainties are derived based on the cumulative histogram of each parameter. For the stellar mass and radius uncertainties, the reported scatter in Torres et al. (2010) is also added in quadrature. Each parameter agrees to within 1σ of the spectroscopic/SED/catalog values, and these parameters are tabulated in Table 3.

### 3.3. Mass of the Candidate Low-mass Companion

Using the spectroscopic orbital elements from the RV fit, we can derive the mass function of the companion,

\[
M_f = \frac{(M_* \sin i)^3}{(M_* + M_c)^2} = \frac{K^2(1-e^2)^{3/2}P}{2\pi G},
\]

\(i\) is the orbit inclination angle.
which is independent of the mass of the primary and the inclination of orbit. For MARVELS-5b, we obtain

$$M_f = (1.742 \pm 0.026) \times 10^{-4} \, M_\odot,$$

where the uncertainty is essentially dominated by the uncertainty in $K$ (see Table 2). Assuming $\sin i = 1$, we derive its minimum mass $M_{\text{min}} = 65.0 \pm 2.9 M_{\text{Jup}}$. The uncertainty here is dominated by the uncertainty in the primary mass (see Table 3). We also find the minimum mass ratio of the companion $q_{\text{min}} = 0.0560 \pm 0.0015$.

The true mass of the companion depends on the inclination of its orbit, which is unknown. We can estimate the posterior probability distribution of the true mass, assuming an isotropic distribution of orbits and adopting a prior for the distribution of the companion mass ratios. We therefore consider three reasonable priors on the companion mass ratio of the form: $dN/dq \propto q^\alpha$, where $\alpha = -1, 0, +1$ (e.g., Grether & Lineweaver 2006). The estimation was realized by using an MCMC, which has been described in detail in Fleming et al. (2010) and Lee et al. (2011). All sources of uncertainty from the mass function and the primary mass have been considered appropriately. We draw values of $\cos i$ from a uniform distribution and weight the resulting distribution by $q^{\alpha+1}$ in order to account for the mass ratio prior. For $\alpha > 0$, the a posteriori distribution does not converge. However, we can rule out mass ratios $q > 1$ for main-sequence companions by the lack of a second set of spectral lines in the high-resolution spectra. We therefore enforce $q \leq 1$, thus implicitly assuming the companion is not a stellar remnant.

The resultant cumulative distributions of the true companion mass are presented in Figure 5, and we summarize the median mass as well as the transit probability for each of our priors in Table 4. For $\alpha < 0$, MARVELS-5b is more likely to be a true BD; for $\alpha = 0$ or $\alpha = 1$, it is more likely to be a low-mass stellar companion.

### 3.4. Evolutionary State of the Host Star

We estimate the evolutionary state of the host star HIP 67526 by comparing the measured stellar parameters with a Yonsei–Yale stellar evolutionary track (Demarque et al. 2004) for an analogous star with $M_\star = 1.10 \, M_\odot$ and [Fe/H] = +0.04. The shaded region indicates the $1\sigma$ deviations in the evolutionary track. The blue dots are the location of the analogous star at different ages in Gyr. HIP 67526 (in red) is most likely a main-sequence dwarf star younger than $\sim 2.5$ Gyr, judging by the evolutionary data alone, since most of the area within the $1\sigma$ ellipsoid lies close to the ZAMS, but its low level of activity suggests an age over $\sim 3$ Gyr, and thus it is most likely a middle-aged star.

(A color version of this figure is available in the online journal.)
3.5. Direct Imaging Search for Visual Companions

3.5.1. FastCam Lucky Imaging

Lucky imaging (LI, observations taken at very high cadence to achieve nearly diffraction-limited images from a subsample of the total) was performed using FastCam (Oscoz et al. 2008) on the 1.5 m Carlos Sánchez Telescope (TCS) at Observatorio del Teide in Spain. The primary goal of these observations was to search for companions at large separations that could contaminate spectroscopic observations of the target masquerading as a systematic trend in the RV data (Fleming et al. 2012). The LI frames were acquired on 2011 April 3, 2011 May 5, and 2011 May 8 in the \( I \) band and spanning \( \sim 21'' \times 21'' \) on the sky. On 2011 April 3 a total of 100,000 short-exposure images, each corresponding to 35 ms exposure time were acquired, on 2011 May 5 a total of 45,000 short-exposure images, each corresponding to 35 ms exposure time were acquired, and on 2011 May 8 a total of 45,000 short-exposure images, each corresponding to 50 ms exposure time were acquired. The data were processed using a custom IDL software pipeline. After identifying frames corrupted due to cosmic rays, electronic glitches, etc., the remaining frames were bias corrected and flat fielded.

LI selection was applied using a variety of selection thresholds (best \( X\% \)) based on the brightest pixel (BP) method. The selected BP must be below a specified brightness threshold to avoid selecting cosmic rays or other non-speckle features. As a further check, the BP must be consistent with the expected energy distribution from a diffraction speckle under the assumption of a diffraction-limited point-spread function (PSF). The BPs of each frame are then sorted from brightest to faintest, and the best \( X\% \) are then shifted and added to generate a final image. In Figure 7, we show the results of the LI selection and shift-and-add for different LI thresholds ranging from considering only the best 1\% of the frames up to including 80\% of the data for data collected in 2011 April and 2011 May. Each panel covers \( \sim 5''.5 \times 5''.5 \) centered on HIP 67526. Restricting the LI selection to the top percentage (i.e., the 1\% LI image) improves the angular resolution with respect to choosing a lower threshold (i.e., the 80\% LI image) but at the cost of higher noise at large distances from the target.

We follow the same procedure as in Femenía et al. (2011) to compute the 3\( \sigma \) detectability (\( \delta m \)) curves on each of the images whose \( \sim 5''.5 \times 5''.5 \) region around HIP 67526 has been depicted in Figure 7: at a given angular distance \( \rho \) from HIP 67526 we identify all possible sets of small boxes of a size larger but comparable to the FWHM of the PSF (i.e., \( 5 \times 5 \) pixel boxes). Only regions of the image showing structures easily recognizable as spikes due to diffraction of the telescope spider and/or artifacts on the read-out of the detector are dismissed. For each of the valid boxes on the arc at angular distance \( \rho \) the standard deviation of the image pixels within the \( 5 \times 5 \) pixel boxes is computed. The value assigned to the 3\( \sigma \) detectability curve at \( \rho \) is three times the mean value from the standard deviations of all the eligible boxes at \( \rho \). This procedure on each of the LI % thresholding values (in steps of 1%) produces a detectability curve, while the envelope of the entire family of curves for a given night yields the best possible detectability curve to be extracted from the whole data set. These “best LI curves” for each of the three nights are depicted in Figure 8, where we can see the data collected are similar in quality to the data on May 8 providing slightly better contrast values.

Figure 7. Composite image showing the results of different LI thresholding on the frames acquired with FastCam at the TCS telescope on 2011 April 3, 2011 May 5 and 2011 May 8. This set of images (in logarithmic scale) illustrates the gain in angular resolution close to the target location when applying high restrictive LI thresholds but at the cost of lowering the contrast achieved at large angular distances from target location (see also Figure 8).

(A color version of this figure is available in the online journal.)

No stellar tertiary to HIP 67526 is detected above the “best LI curves.”

3.5.2. Keck Adaptive Optics Imaging

To further assess the multiplicity of HIP 67526, we acquired high angular resolution images of the star on UT 2012 24 June using NIRC2 (instrument Pf. K. Matthew) with the Keck II AO system (Wizinowich et al. 2000). AO observations probe the immediate vicinity of host stars and generate deep contrast compared to LI (e.g., Fleming et al. 2012; Ma et al. 2013). Furthermore, AO observations are sensitive to objects with red colors given the nominal 1–3 \( \mu m \) wavelength operating range.

Our observations consist of dithered frames taken with the \( K' \) (\( \lambda_c = 2.12 \mu m \)) filter. We used the narrow camera setting to provide fine spatial sampling of the NIRC2 PSF. The total on-source integration time was 190 s. Images were processed using standard techniques to replace hot pixel values, flat field the detector array, subtract thermal background noise, and align and coadd frames.
Keck AO image of HIP 67526. No stellar companions are detected with $\Delta m_K < 5$ mag for separations beyond 0′25 and $\Delta m_K < 8$ mag for separations beyond 1′0 at 10σ significance level (see also Figure 10).

(A color version of this figure is available in the online journal.)

Figures 9 and 10 show the final reduced AO image and corresponding contrast curve. No candidate companions were noticed in individual raw frames or the final reduced image. Our diffraction-limited observations rule out the presence of companions with $\Delta m_K < 5$ mag for separations beyond 0′25 and $\Delta m_K < 8$ mag for separations beyond 1′0 (10σ). We employ the empirical mass–luminosity relationships in Delfossé et al. (2000) to derive the upper mass limit of the undetected companions; this analysis results in an upper mass limit 0.2 $M_\odot$ for separations larger than 40 AU and 0.1 $M_\odot$ for separations larger than 100 AU.

4. DISCUSSION AND SUMMARY

The frequency of BD companions to solar-like stars at close and intermediate separations is less than 1% (Marcy & Butler 2000), which is much less than the frequency of planetary companions (>10%; e.g., Howard et al. 2010; Mayor et al. 2011) and the frequency of spectroscopic stellar binaries detected in RV surveys (~14%; e.g., Halbwachs et al. 2003). The frequency of BD companions was recently updated by Sahlmann et al. (2011) to be <0.6% on the basis of the CORALIE planet-search sample. This result is more accurate since the authors ruled out companions having true masses in the stellar regime using the Hipparcos astrometric measurements to determine the orbital inclinations. Constraining the mass distribution of companions can provide an important observational clue to distinguish the formation and evolution mechanism of planetary, BD, and stellar companions. The current mass distribution suggests that low-mass BD companions less than $\sim 30 M_{\text{Jup}}$ are likely to form in protoplanetary disks, while companions more massive than $\sim 45 M_{\text{Jup}}$ form via fragmentation (Grether & Lineweaver 2006; Sahlmann et al. 2011; Ma & Ge 2013). The BD and low-mass stellar companion discoveries from MARVELS will result in a more precise determination of the mass limits of core accretion and gravitational collapse. MARVELS-5b contributes to constraining the shape of the massive BD–low-mass star boundary.

Spectroscopic binaries generally show moderately eccentric orbits (e.g., Duquennoy & Mayor 1991; Raghavan et al. 2010). Ribas & Miralda-Escudé (2007) reported a tentative trend that low-mass planets ($M \sin i < 4 M_{\text{Jup}}$) generally have lower eccentricity than high-mass planets ($M \sin i > 4 M_{\text{Jup}}$), having a similar eccentricity distribution as binary stars (Figure 3 of Ribas & Miralda-Escudé 2007). Díaz et al. (2012) reported that most of the BD companions in their sample exhibit a considerable orbital eccentricity, supporting the eccentricity-mass trend. MARVELS-5b has a high eccentricity (0.44), which is around the peak of the eccentricity distribution of the observed BD and low-mass stellar companions (Sahlmann et al. 2011; Díaz et al. 2012). In view of MARVELS-5b’s eccentricity, it is probably a member of the main population of these massive companions to solar-like stars. Our previous MARVELS discoveries (MARVELS-2,3,4,6b) all have an eccentricity lower than ~0.2 (De Lee et al. 2013).

A stellar tertiary is likely to affect the formation and evolution of the substellar companion to the primary. Observationally, Zucker & Mazeh (2002) point out that planets found in binaries may have a negative period–mass correlation rather than the positive correlation between the masses and periods of the...
planets orbiting single stars. By studying a larger sample (19 planets in a double or multiple star system), Eggenberger et al. (2004, 2007) showed that short-period \( (P < 40 \text{ days}) \) planets found in multiple star systems may follow a different period–eccentricity distribution than the short-period planets around isolated stars. These observations seem to indicate that the presence of a stellar companion alters the migration and mass growth rates of planets (Kley 2001). Similar influences have also been observed on close spectroscopic binaries in triple systems. Shorter period binaries are more likely to be in multiple-star systems, i.e., \( \sim 80\% \) for \( P < 7 \text{ days} \) versus \( \sim 40\% \) for \( P > 7 \text{ days} \) (Tokovinin et al. 2006). This significant difference suggests that the periods of close binary systems with triples were efficiently decreased by angular momentum exchange with companions.

With masses between planetary companions and stellar components in spectroscopic binaries, the formation and migration of BD and low-mass stellar companions can certainly be affected by the presence of a tertiary as well. However, this problem has not been studied in a statistical way since the current BD and low-mass stellar companion sample is fairly small and no systematic survey of stellar tertiaries for these companions has been conducted. Using high-contrast imaging, the MARVELS survey goes to great lengths to investigate the statistics of its own discoveries of low-mass companions in the presence/absence of a stellar tertiary. As mentioned in Section 1, most of the previous MARVELS discoveries have a stellar tertiary (or a candidate stellar tertiary) detected either by high-contrast imaging or analysis of the long-term RV trend. Among the confirmed discoveries, MARVELS-3b (Wisniewski et al. 2012) has an orbital period \( (P \sim 79 \text{ days}) \) and minimum mass ratio \( (q_{\text{min}} \sim 0.09) \) similar to MARVELS-5b (this work), but the former has a less eccentric orbit \( (e \sim 0.1) \). Wisniewski et al. (2012) found a faint candidate tertiary companion on the Keck AO image, separated by \( \sim 1'' \) from the primary, thus speculating that MARVELS-3b might have initially formed in a tertiary system with much different orbital parameters and reached its current short-period orbit during the cluster dispersal phase. For MARVELS-5b, the Keck AO imaging rules out any star with mass greater than \( 0.1 \ M_\odot \) at a separation larger than \( 1'' \) from the primary. This may imply that other formation mechanisms of low-mass-ratio binaries are needed.

In summary, we report a candidate BD or low-mass stellar companion to the solar-like star HIP 67526. The best Keplerian orbital fit parameters were found to have an orbital period of \( 90.2695^{+0.0188}_{-0.0187} \) days, an eccentricity of \( 0.4375 \pm 0.0040 \), and a semi-amplitude of \( 2948.14^{+16.65}_{-16.55} \text{ m s}^{-1} \). The minimum companion mass was determined to be \( 65.0 \pm 2.9 M_\text{Jup} \). This object helps to populate the high-mass end of the sparsely populated region of the mass function of companions to solar-type stars and provide observational evidence to constrain formation and evolution theories. No stellar tertiary is detected with high-contrast imaging for the MARVELS-5 system, while all the other previous MARVELS-discovered systems appear to have at least one stellar companion.

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